THE OCEAN BOMB RADIOCARBON INVENTORY REVISITED

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ABSTRACT. Large discrepancies exist among data-based estimates and model reconstructions of the ocean bomb radiocarbon inventory. In order to resolve this gap, it has been proposed that the CO_2 piston velocity should be revised downward (Sweeney et al. 2007; Müller et al. 2008). This article compares the transient ^{14}C distributions in the ocean obtained with different formulations of the isotopic ratio commonly used in modeling studies. It is found that both the CO_2 increase and the air-sea CO_2 flux significantly contribute to the 1990 ocean bomb ^{14}C inventory, by around 10% each. Moreover, these 2 processes explain more than 25% of the inventory difference between 1974 and 1990. These results imply that, as already argued by Naegler (2009), inventories based on observations that lack information about CO_2 invasion are underestimated. Further, this work provides insight into the reasons for discrepancies among model results. It suggests that while a comprehensive isotopic formulation is needed when addressing the global ^{14}C cycle, a simplified form is more relevant for model calibration and piston velocity assessment based on currently available bomb ^{14}C inventories.

INTRODUCTION

Our ability to predict future Earth climate relies on our understanding of the global carbon cycle as well as of the ocean circulation. Ocean ventilation, i.e. the replacement of interior waters by seawater having been in contact with the atmosphere, is a key climatic process since it controls the rate at which the ocean can sequester excess heat and carbon. On the decadal to centennial timescales, the magnitude of the atmospheric CO_2 disturbance is determined by exchange with the terrestrial and surface ocean reservoirs. The radiocarbon released by atmospheric nuclear weapons testing ("bomb" ^{14}C) is an ideal tracer to gain insight into the specific rates characterizing the carbon cycle and ocean ventilations on such timescales. Assessment of the global distribution of bomb ^{14}C provides constraints on the carbon fluxes between reservoirs (e.g. Hesshaimer et al. 1994; Lassey et al. 1996; Levin et al. 2010). More specifically, it allows determination of the global air-sea gas exchange rate (Wanninkhof 1992; Krakauer et al. 2006; Sweeney et al. 2007; Naegler 2009) and therefore to estimate the net air-sea CO_2 flux (Takahashi et al. 2009). It also offers the opportunity to assess near-surface circulation, vertical transport, and mixing schemes in OGCMs (Toggweiler et al. 1989b; Jain et al. 1995; Joos et al. 1997; England and Rahmstorf 1999; Orr et al. 2001; Müller et al. 2006).

For these reasons, a precise knowledge of the ocean bomb ^{14}C inventory is needed. However, this quantity has been long debated and seems elusive. The initial estimate by Broecker et al. (1985) has either been questioned (Hesshaimer et al. 1994; Duffy and Caldeira 1995) or corroborated (Broecker et al. 1995; Lassey et al. 1996). Since then, new methods (e.g. Rubin and Key 2002; Peacock 2004) and field-based measurements (Key et al. 2004) have been made available. With the exception of an increased inventory in the Southern Ocean (Leboucher et al. 1999), recent works based on field studies usually call for a downward revision of the bomb ^{14}C inventory (e.g. Peacock 2004; Sweeney et al. 2007). An important implication of these results is that the air-sea gas exchange rate initially proposed by Wanninkhof (1992) should be significantly reduced: Sweeney et al. (2007) and Müller et al. (2008) suggest reductions by as much as 30% and 19%, respectively.

The use of bomb ^{14}C as an ocean tracer is not straightforward because the exact amount that entered the sea may only be indirectly obtained. Indeed, except for measurements on corals or seashells, we do not have any direct evidence of pristine ocean values (Broecker et al. 1995; Rubin and Key

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2002). Bomb ^{14}C is computed by subtracting the observed (post-bomb) ^{14}C concentrations from an estimate of the natural pre-anthropogenic values (Broecker et al. 1995). The natural ^{14}C is reconstructed on the basis of empirical relationships between $\Delta^{14}C$ and other tracers (Broecker et al. 1995; Rubin and Key 2002; Peacock 2004; Sweeney et al. 2007). These procedures also necessitate the separation of bomb-free water masses from those already contaminated. This task is performed with the help of transient tracers such as CFCs or tritium or with the help of OGCMs. Other factors significantly contribute to the uncertainty of field-based estimates. First, incomplete spatial and temporal coverage as well as significant delay between data collection in different ocean basins produce biased estimates. The second significant source of uncertainty comes from the fact that bomb inventories are computed on the basis of available DIC samples contemporaneous to $\Delta^{14}C$ measurements (Key et al. 2004). Hence, data-based inventories probably provide a lower-bound estimate because they neglect the fact that DIC levels were lower in earlier decades (Naegler 2009).

Model studies, on the other hand, do not yet provide a means for solving these difficulties. Bomb ^{14}C ocean inventories as estimated by different models encompass a wide range of values (Joos 1994; Duffy and Caldeira 1995; Joos and Bruno 1998; Müller et al. 2008) and are usually much larger than data-based inventories. Processes controlling the bomb ^{14}C uptake in the ocean include the air-sea exchange rate but also the renewal rate of surface waters (the ventilation), which is model-dependent. Joos et al. (1997), Joos and Bruno (1998), and Müller et al. (2008) attribute most of the discrepancy to the gas exchange coefficient. Other sources for differences in model predicted values include surface boundary conditions in the Southern Ocean (Toggweiler and Samuels 1993), vertical mixing (Toggweiler et al. 1989b; Müller et al. 2008), or the numerical scheme (Bryan 1987; Duffy et al. 1997).

There is another source of discrepancy among model results that has been overlooked. It is to be found in the way ^{14}C is represented in models (models that sometimes will complement field information to provide data-based estimates; e.g. Naegler et al. 2006; Sweeney et al. 2007). Several model studies adopted the simplified formulation of Toggweiler et al. (1989a) to describe the transport of ^{14}C in the ocean (Maier-Reimer et al. 1993; Stocker et al. 1994; Duffy and Caldeira 1995; Campin et al. 1999; England and Rahmstorf 1999; Rodgers et al. 2000; Mahadevan 2001; Orr et al. 2001; Butzin et al. 2005; Müller et al. 2006). This formulation does not much affect the equilibrium (or pre-industrial) ^{14}C distribution in the ocean (Maier-Reimer 1993; Orr et al. 2001). However, its effect on the ocean bomb ^{14}C uptake might be more significant since it relies on the hypothesis that air-sea CO_2 disequilibrium remains constant over the period of interest.

We examine the transient distributions of ^{14}C in the ocean from modeling studies based on different formulations for ^{14}C. For this purpose, we use a 3D model of the carbon cycle driven by 2 different circulation fields. Two methods for representing ^{14}C are then implemented in the model. With one formulation all carbon isotopes are transported while with the other $\Delta^{14}C$ is considered as the state variable. The differences between the results obtained with the 2 methods as well as the implications for estimates of the CO_2 gas exchange coefficient based on ^{14}C bomb inventories are examined and discussed. Methods and model setup are presented below. The description of numerical experiments results precedes the general discussion and conclusions.

METHODS

Modeling Radiocarbon

A distinction is made between isotopic ratios relative to the number of ^{12}C atoms in the sample ($\Re$) or to the total number of carbon atoms (R). Hence, we define the ratios $^{13}R = {}^{13}C/C$, $^{14}R = {}^{14}C/C$, and

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$^{13}\Re = {}^{13}C/{}^{12}C$. R is the ratio most readily computed in numerical models while field measurements are usually expressed using the following quantities:

$$\delta^{13}C = \left(\frac{^{13}\Re}{^{13}\Re_{PDB}} - 1\right)10^3 = \left(\frac{^{13}R}{(1 - {}^{13}R)\,^{13}\Re_{PDB}} - 1\right)10^3 \tag{1}$$

$$\Delta^{14}C = \left(\frac{^{14}R}{^{14}R_{OXA}}\left(\frac{0.975}{1 + \delta^{13}C/10^3}\right)^2 - 1\right)10^3 \tag{2}$$

in which $^{13}\Re_{PDB} = 1.12372 \times 10^{-2}$ is the PDB standard ratio, $^{14}R_{OXA} = 1.176 \times 10^{-12}$ the oxalic acid standard, and 0.975 a constant term corresponding to the average terrestrial wood $^{13}\Re$ ratio for the pre-industrial period (Stuiver and Polach 1977; Keeling 1981).

With **u** the 3-D velocity field and *K* the diffusivity tensor, the equations governing the transport of total dissolved inorganic carbon concentration C_T and that of its isotopes ^iC_T (with i = 13 or 14) in the ocean read

$$\frac{\partial}{\partial t}C_T = -\nabla \cdot (\mathbf{u}C_T - K \cdot \nabla C_T) + q_v(C_T) \tag{3}$$

$$\frac{\partial}{\partial t}{}^{13}C_T = -\nabla \cdot (\mathbf{u}\,{}^{13}C_T - K \cdot \nabla\,{}^{13}C_T) + q_v({}^{13}C_T) \tag{4}$$

$$\frac{\partial}{\partial t}{}^{14}C_T = -\nabla \cdot (\mathbf{u}\,{}^{14}C_T - K \cdot \nabla\,{}^{14}C_T) + q_v({}^{14}C_T) - \lambda\,{}^{14}C_T \tag{5}$$

with λ the decay rate for ^{14}C. The term $q_v(X)$ represents appropriate source minus sink term for tracer X. It takes into account biotic or abiotic (e.g. $CaCO_3$ dissolution) processes, at the exclusion of radioactive decay.

At the ocean-atmosphere interface, we impose a Neumann boundary condition. With $\Im$ the carbon flux through the air-sea interface, the boundary condition for C_T reads

$$\mathrm{n} \cdot (\mathrm{K} \cdot \nabla \mathrm{C_T}) = \Im = \mathrm{F^{as}} - \mathrm{F^{sa}} \tag{6}$$

where n is a unit vector normal to the interface. The individual flux components are given by

$$\mathrm{F^{as}} = \kappa_{CO2}\mathrm{K_H}\mathrm{p^a_{CO2}} \text{ and } \mathrm{F^{sa}} = \kappa_{CO2}\mathrm{CO_{2s}} \tag{7}$$

where the superscripts *as* and *sa* refer to the air-to-sea and sea-to-air transfers, respectively. In (7), κ_{CO2} is the carbon dioxide transfer velocity, K_H the CO_2 solubility in seawater, p^a_{CO2} the atmospheric pressure at sea level, and CO_{2s} the concentration of dissolved carbon dioxide in seawater, which is obtained by computing carbonate speciation.

The air-sea boundary conditions for isotopes are based on (6) and (7) with fractionation factors and isotopic ratios weighting the different fluxes (Heimann and Maier-Reimer 1996). These fluxes read

$$^{13}\Im = {}^{13}\alpha^{as}\,{}^{13}R_a F^{as} - {}^{13}\alpha^{sa}\,{}^{13}R_{C_T} F^{sa} \tag{8}$$

$$^{14}\Im = {}^{14}\alpha^{as}\,{}^{14}R_a F^{as} - {}^{14}\alpha^{sa}\,{}^{14}R_{C_T} F^{sa} \tag{9}$$

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Here, ${}^{i}R_a$ and ${}^{i}R_{C_T}$ represent isotopic ratios (i = 13 or 14) in atmospheric CO_2 and in ocean DIC, respectively. Fractionation effects between atmospheric and dissolved carbon dioxide are represented by ${}^{13}\alpha^{as}$ while ${}^{13}\alpha^{sa}$ contains the fractionation factors between dissolved CO_2 and carbonate species.

The equations (3)–(5) together with the boundary conditions (6)–(9) and appropriate expressions for the source and sink terms constitute the complete set of equations of use in ^{14}C modeling. In the following, we refer to this calculation as the F method.

A Simplified Formulation

Since $\Delta^{14}C$ represents the ratio that would have been observed if isotopic fractionation did not occur, it may be used as such as a physical tracer of ocean ventilation (Broecker et al. 1961; Maier-Reimer 1993). Hence, as was initially suggested by Toggweiler et al. (1989a,b), in many model studies the ratio ^{14}R is transported rather than the individual concentrations. By combining the transport equations (3) and (5), one obtains the equation governing the evolution of the ratio ${}^{14}R_{C_T}$

$$\frac{\partial}{\partial t}{}^{14}R_{C_T} = -\nabla\cdot(u\,{}^{14}R_{C_T} - K\cdot\nabla\,{}^{14}R_{C_T}) + \frac{1}{C_T}(q_v({}^{14}C) - {}^{14}R_{C_T}q_v(C_T)) - \lambda\,{}^{14}R_{C_T} + \frac{2}{C_T}\nabla C_T\cdot K\cdot\nabla\,{}^{14}R_{C_T} \quad (10)$$

The last term on the right-hand side is generated by mixing processes and has no equivalent in the equation governing concentrations. This term may, in most parts of the deep ocean, be considered as a sink term. Indeed, the oldest ocean water (lowest ^{14}R) would generally correspond to those with the highest C_T content.

Solving Equation 10 would call for a specific numerical scheme, different from the one applied to the other variables in the model (temperature, salinity, tracers...) with consequences that ^{14}R would loose its relevance for diagnosing the model circulation and ventilation timescale. Therefore, in order to represent the ratio ^{14}R in the numerical model, the mixed product in Equation 10 must vanish. It is equivalent to assuming that the C_T field is homogeneous and constant, i.e.

$$C_T(x, t) = \overline{C_T} \quad (11)$$

with $\overline{C_T}$ a typical DIC concentration and x the position. Further, since ^{14}R is considered as a purely physical tracer, biological activity and fractionation factors are ignored. Under these hypotheses, (10) becomes

$$\frac{\partial}{\partial t}{}^{14}R_{C_T} = -\nabla\cdot(u\,{}^{14}R_{C_T} - K\cdot\nabla\,{}^{14}R_{C_T}) - \lambda\,{}^{14}R_{C_T} \quad (12)$$

Denoting ${}^{R}\Im$ the flux of ^{14}R through the air-sea interface, the boundary condition to be applied to (12) is

$$n\cdot(K\cdot\nabla\,{}^{14}R_{C_T}) = {}^{R}\Im \quad (13)$$

This condition may be expressed in terms of the boundary conditions for C_T and ^{14}C. Indeed, with the help of Equations 6, 7, and 9 in which we neglect fractionation, (13) becomes:

$${}^{R}\Im = n\cdot(K\cdot\nabla\,{}^{14}R_{C_T}) = n\cdot\left(K\cdot\left(\frac{\nabla\,{}^{14}C_T}{C_T} - {}^{14}R_{C_T}\frac{\nabla C_T}{C_T}\right)\right)$$

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$$= \frac{1}{C_T}\left({}^{14}\mathfrak{I} - {}^{14}R_{C_T}\mathfrak{I}\right)$$

$$= \frac{1}{C_T}[({}^{14}R_a F^{as} - {}^{14}R_{C_T}F^{sa}) - {}^{14}R_{C_T}(F^{as} - F^{sa})] \tag{14}$$

Then, with assumption (11) the flux ${}^{R}\mathfrak{I}$ reads

$${}^{R}\mathfrak{I} = \frac{\kappa_{CO_2}K_H}{\overline{C_T}}p^a_{CO_2}({}^{14}R_a - {}^{14}R_{C_T}) \tag{15}$$

The implications of Equation 15 may be understood by comparing it to (9) which, with the help of (7) and ignoring the minor effect of fractionation, reads

$${}^{14}\mathfrak{I} = \kappa_{CO_2}({}^{14}R_a K_H p^a_{CO_2} - {}^{14}R_{C_T}CO_{2s}) = \kappa_{CO_2}K_H\left({}^{14}R_a p^a_{CO_2} - {}^{14}R_{C_T}p^s_{CO_2}\right)$$

$$= \kappa_{CO_2}K_H(p^a_{CO_2}({}^{14}R_a - {}^{14}R_{C_T}) + (p^a_{CO_2} - p^s_{CO_2})^{14}R_{C_T}) \tag{16}$$

where $p^s_{CO_2}$ is the partial pressure of CO_2 in the ocean. The second term on the right-hand side of (16) represents the contribution of air-sea CO_2 disequilibrium to the ^{14}C flux. Since this term does not intervene in the computation of ${}^{R}\mathfrak{I}$, the form (15) is equivalent to assuming local air-sea CO_2 equilibrium. It means that any method based on (15) postulates that the ^{14}C fluxes into the ocean have a time history directly proportional to the atmospheric ^{14}C perturbation and ignores the role of air-sea CO_2 flux in setting the ^{14}C exchange rate. The method represented by the set of Equations 12 and 15 will hereafter be referred to as the R method.

Model Setup

We use the ocean carbon cycle model LOCH (Mouchet and François 1996; Menviel et al. 2008; Goosse et al. 2010) in its abiotic version to simulate ocean carbon chemistry. We restrict the present discussion to abiotic processes since it may be shown that the impact of biological activity on the global figures discussed here is negligible (Müller et al. 2008; Mouchet 2011). LOCH is driven off-line by three-dimensional fields of velocity components, potential temperature, salinity, and convective mixing from 2 OGCMs in their annually averaged versions: the Large-Scale Geostrophic Ocean General Circulation Model (LSG-OM; Maier-Reimer et al. 1993; Mikolajewicz et al. 1993) and the Louvain-la-Neuve OGCM (UL-OM; Deleersnijder and Campin 1995; Campin and Goosse 1999). These models belong to the coarse-grid OGCM class but nevertheless capture the key features of ocean circulation. The use of both circulation fields is motivated by the differences they exhibit in the processes controlling the distribution of tracers and in the circulation characteristic scales (Mouchet et al. 2012). Diffusion in UL-OM is mostly explicit with maxima in the upper layers and in convective areas. This is in contrast to LSG-OM in which numerical diffusion (a more ubiquitous process) prevails in setting the effective diffusivity. For a discussion of these processes and more details on the OGCMs used in this study, see Mouchet et al. (2012) and references therein.

The gas transfer velocity is based on the empirical formulation of Wanninkhof (1992). We consider a reduced air-sea exchange rate in the presence of sea ice. Therefore, the piston velocity in the model is computed in the following way

$$\kappa_{CO_2} = (1 - f_{ice})K_W\overline{w^2}\sqrt{660/Sc} \tag{17}$$

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with f_{ice} the fraction of the sea surface covered with ice and *Sc* the Schmidt number. The long-term average of the squared wind magnitude $\overline{w^2}$ is taken from the annual climatology of Trenberth et al. (1989). The annual average of sea-ice fraction is computed with the help of the OCMIP-2 sea-ice cover climatology (Orr et al. 2000). Both wind and sea ice fields were carefully interpolated onto both OGCM grids to reduce the differences between carbon cycle experiments to a strict minimum.

The Schmidt number *Sc* is calculated with the help of the formulation of Wanninkhof (1992). The CO_2 solubility (taken from Weiss 1974) as well as *Sc* are computed with the OGCM temperature and salinity fields. The coefficient K_W in (17) is then adjusted to obtain an average value of the CO_2 invasion rate (i.e. $\kappa_{CO2} \times$ solubility) of the order of 0.061 mol ppm^{-1} m^{-2} yr^{-1} with an atmospheric CO_2 of 280 ppm (Siegenthaler 1986; Watson and Liss 1998). In the reference experiments, the globally averaged CO_2 piston velocity is 17.2 cm/hr with LSG-OM and 17.1 cm/hr with UL-OM.

The temperature-dependent formulations for the fractionation factors in (8) are from Vogel et al. (1970) and Mook et al. (1974), and the small kinetic effect on gas exchange velocities is taken from Siegenthaler and Münnich (1981). Fractionation factors for ^{14}C are the squares of those for ^{13}C (Siegenthaler and Münnich 1981). Since the model considers the full carbonate chemistry and since individual fractionation factors are applied between the different dissolved carbonate species, the total fractionation occurring during air-sea exchange also depends on CO_2 concentration.

We performed experiments over the industrial era in which the ocean ^{14}C distribution is computed according to the F or R method. In addition, 2 experiments address the sensitivity of bomb ^{14}C inventories to the transfer velocity; in these K_W is increased by 10% and decreased by ~15%, respectively. Each of the ^{14}C modeling methods (F or R) requires its own initial conditions. This rule also prevails for experiments addressing the sensitivity to piston velocity. In total, 4 pre-industrial states were prepared with each circulation fields in order to provide the initial conditions to the corresponding transient experiments (Table 1). The $\Delta^{14}C$ values in the ocean are initialized to a globally uniform pre-bomb value of −140‰. The model is then run to equilibrium constrained by pre-industrial atmospheric values of CO_2, $\delta^{13}C$, and $\Delta^{14}C$ (277.95 ppm, −6.51‰, and 0‰, respectively). For the value of $\overline{C_T}$ intervening in Equation 15, we use the globally averaged surface C_T concentration obtained from the equilibrium experiment F0. To simulate ocean ^{14}C during the industrial era, we force atmospheric CO_2, $\delta^{13}C$, and $\Delta^{14}C$ to follow the historical estimates. CO_2 and $\Delta^{14}C$ are from Enting et al. (1994) and OCMIP-2 (Orr et al. 2000) while $\delta^{13}C$ data are from Francey et al. (1999) up to 1978 and from Keeling et al. (2010) onward. Two additional experiments are performed in which the atmospheric CO_2 (and consequently ^{13}R) is maintained at its pre-industrial level. The purpose of these additional experiments is to illustrate the role of the anthropogenic CO_2 increase in the bomb ^{14}C ocean uptake. Transient experiments are identified this way: the first letter refers to the ^{14}C modeling method (R or F), the second letter to the type of CO_2 forcing over the industrial era (C: CO_2 maintained at pre-industrial value, or P: CO_2 from the historical record), and the third letter, when present, to the magnitude of κCO_2 (L or H, respectively) with respect to the reference state. Hence, 6 transient experiments were carried on with each circulation fields. They are summarized in Table 1. Except for FC, each one of these methods has its equivalent in the literature. The analysis of Joos et al. (1997) and Müller et al. (2008) are performed in a similar way as FP. The F method with biological cycle is considered in the works of Bacastow and Maier-Reimer (1990) and Heimann and Maier-Reimer (1996). The transient experiments of Toggweiler et al. (1989b), Rodgers et al. (2000), Mahadevan (2001), and Ito et al. (2004) are identical to those that we call RC. The ^{14}C experiments during OCMIP-1 (Orr et al. 2001) and in Müller et al. (2006) are based on the RP protocol. In addition, Krakauer et al. (2006) and Sweeney et al. (2007) rely at some stage for their field estimates on model experiments performed the same way as RP.

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Table 1 Summary of the modeling experiments setup. Experiment name and method are given in columns 1 and 2, with the initial condition (IC) for transient experiments in column 3. Column 4 contains the multiplicative factor of the reference CO_2 transfer velocity. Columns 5 to 7 display the atmospheric forcing (pre-industrial PI or industrial I) for CO_2, $\delta^{13}C$, and $\Delta^{14}C$, respectively.

Name	Method	IC	$\times \kappa_{CO2}$	$p^a_{CO_2}$	$^{13}R_a$	$^{14}R_a$
R0	R	—	1	PI	—	PI
F0	F	—	1	PI	PI	PI
F0L	F	—	0.85	PI	PI	PI
F0H	F	—	1.10	PI	PI	PI
RC	R	R0	1	PI	—	I
RP	R	R0	1	I	—	I
FC	F	F0	1	PI	PI	I
FP	F	F0	1	I	I	I
FPL	F	F0L	0.85	I	I	I
FPH	F	F0H	1.10	I	I	I

RESULTS

Pre-Industrial State

As illustrated in Figure 1, there is a fair agreement between model results in the upper ocean and the "natural" or pre-bomb $\Delta^{14}C$ distribution as reconstructed by Key et al. (2004). The contrast between the Southern Ocean and the other ocean basins is reasonably reproduced despite less-negative $\Delta^{14}C$ values in the upper ocean predicted by the models with respect to the GLODAP reconstruction. Some of this discrepancy could result from the Suess effect, which decreased surface ocean $\Delta^{14}C$ by 6 to 12‰ from 1900 to 1952 (Druffel and Suess 1983). The data-based reconstruction meant to be representative of the pre-bomb state (i.e. the early 1950s; Rubin and Key 2002) might indeed be affected to some extent by this process, which is not present in the modeled pre-industrial states. However, model shortcomings cannot be ruled out. Profiles in the deep Atlantic Ocean agree reasonably with data, but both models fail to reproduce deep $\Delta^{14}C$ in the Southern and Pacific Oceans. One may not invoke the neglect of biological processes to explain such departures. Biological activity actually contributes to the ^{14}C pool at depth (Craig 1969), but its impact is nearly totally eliminated by the use of fractionation-corrected ^{14}R (Stuiver et al. 1981; Bacastow and Maier-Reimer 1990; Maier-Reimer 1993; Joos et al. 1997). The main reasons for such discrepancies between model and data are probably to be found in the representation of convective events and vertical mixing processes in coarse grid OGCMs, which generally do not simulate the deep natural ^{14}C distribution well (e.g. Orr et al. 2001; Matsumoto et al. 2004). However, this shortcoming does not necessarily affect processes characterized by short timescales, such as the one at stake.

Pre-industrial or equilibrium $\Delta^{14}C$ distributions do not differ much among the different configurations. The close correspondence of R0 and F0 profiles implies that the contribution of the mixed product term in (10) is very small. The R0 and F0 profiles in the North Atlantic depart more from each other in the LSG-OM case than in the UL-OM case. This is due to the representation of freshwater fluxes in the OGCMs. In LSG-OM, these fluxes are treated as water fluxes and incorporated in the velocity field divergence. In experiments F0, these fluxes affect the individual tracers and mostly cancel out when ratios are computed. In experiments with the R method, these implicit fluxes directly affect the ^{14}R ratio. Evaporation is comparatively high in the North Atlantic Ocean, which leads to increased concentrations in that area. This problem is not present with the UL-OM fields since evaporation and precipitation are represented by means of an explicit flux condition at the

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ocean surface. When performing R experiments with UL-OM, we set these fluxes to zero. We did not attempt to correct for these fluxes in experiments performed with LSG-OM and addressing the R method. Indeed, the results presented hereafter suggest that the influence of these fluxes on estimated inventories is very small with respect to other factors elaborated below.

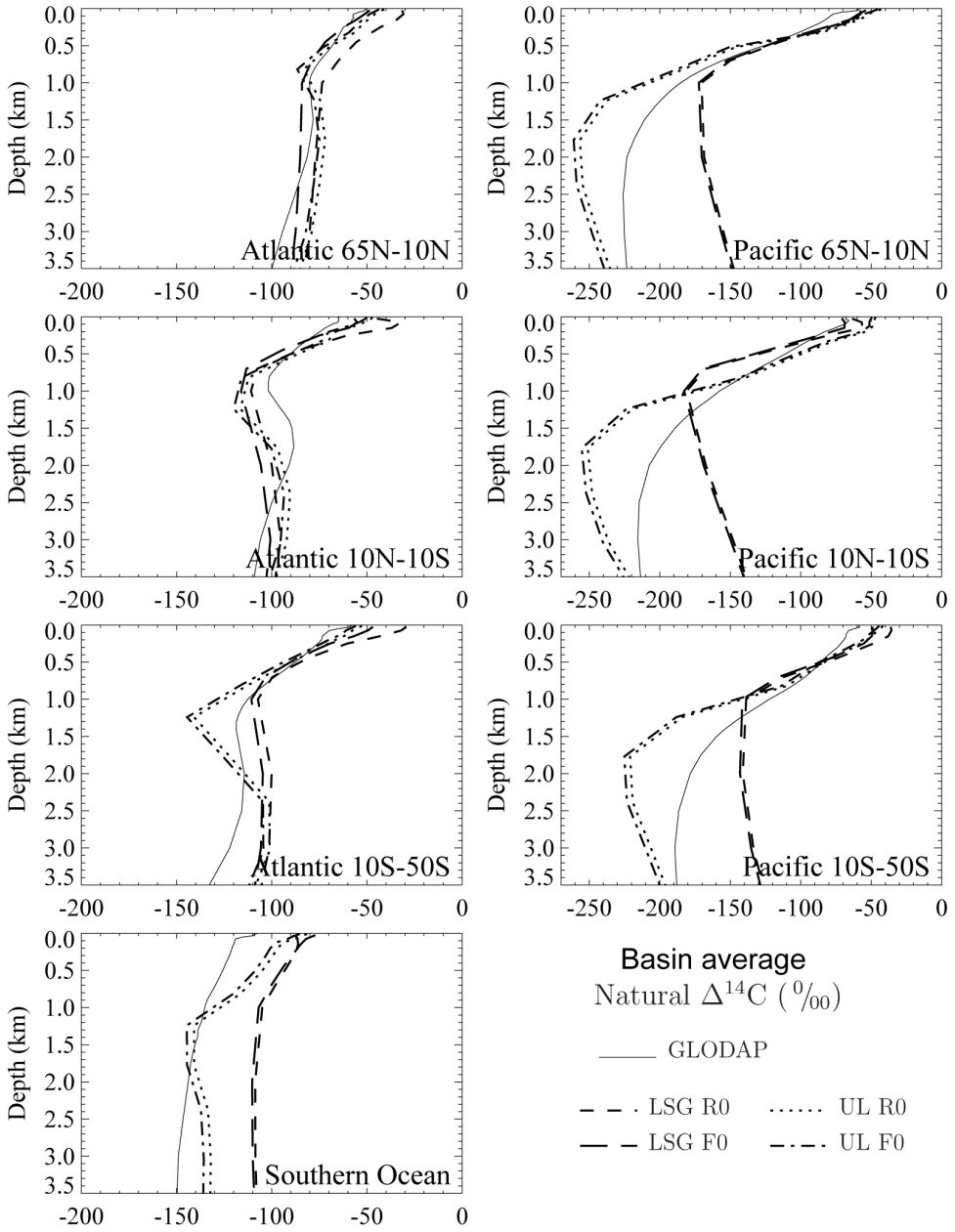


Figure 1 Mean natural $\Delta^{14}C$ vertical profiles in several ocean basins from the GLODAP reconstruction (solid lines; Key et al. 2004) and from experiments F0 and R0 with circulation fields from LSG-OM (LSG) and UL-OM (UL) OGCMs. The values in the north, equatorial, and southern Atlantic and Pacific oceans are reproduced, respectively, in the upper left and right 3 panels, with the bottom panel representing the Southern Ocean (south of 50°S).

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Response to Anthropogenic Forcing

The global net CO_2 fluxes into the ocean during the 1980s and the 1900s amount to 1.86 GtC/yr and 2.15 GtC/yr, respectively, for experiment FP with both circulation fields (Table 2). These values are within the range of estimates based on observations and model studies (Matsumoto et al. 2004; Gruber et al. 2009). The inventory of anthropogenic carbon in the ocean predicted with both circulation fields ($^{C}I_{1800}^{1994}$ in Table 2) is also in good agreement with inventories based on observations (118 ± 19 PgC; Sabine et al. 2004).

Table 2 Simulated ocean bomb ^{14}C inventories and ocean CO_2 uptake. $^{14}I_{1940}^{1975}$ and $^{14}I_{1940}^{1990}$ are the ^{14}C inventories (in 10^{26} at.) relative to 1940 for 1975 and 1990, respectively. $^{14}I_{1975}^{1990}$ is the net bomb ^{14}C uptake from 1975 to 1990 (i.e. the difference between the 2 former inventories). $\overline{\Im}_{1980s}$ and $\overline{\Im}_{1990s}$ correspond to the CO_2 uptake rates (GtC/yr) during the 1980s and 1990s, respectively. The last line contains $^{C}I_{1800}^{1994}$, the cumulative ocean CO_2 uptake (GtC) over the period 1800–1994.

	UL-OM						LSG-OM					
	RC	RP	FC	FP	FPL	FPH	RC	RP	FC	FP	FPL	FPH
$^{14}I_{1940}^{1975}$	239	272	236	288	257	306	244	277	241	294	262	312
$^{14}I_{1940}^{1990}$	313	355	309	388	352	407	322	365	318	399	361	418
$^{14}I_{1975}^{1990}$	74	83	73	100	95	101	78	88	77	105	99	106
$\overline{\Im}_{1980s}$	–	–	–	1.86	1.81	1.89	–	–	–	1.86	1.82	1.88
$\overline{\Im}_{1990s}$	–	–	–	2.15	2.10	2.18	–	–	–	2.15	2.11	2.17
$^{C}I_{1800}^{1994}$	–	–	–	116	119	121	–	–	–	119	121	122

The simulated post-bomb $\Delta^{14}C$ ocean distribution is illustrated in Figure 2. The ^{14}C increase is reasonably well captured by the model. Bomb ^{14}C inventories with UL-OM and with LSG-OM are similar (Table 2), with the latter predicting slightly larger values. On the other hand, the spread among results with different ^{14}C modeling methods is much larger than the differences related to circulation fields. Experiments FC and RC predict similarly low inventories. Larger inventories are obtained with RP experiments while FP predicts the largest uptake.

The low RC and FC inventories are a direct consequence of neglecting the CO_2 increase over the industrial era. Lower CO_2 levels imply lower amounts of ^{14}C atoms for an identical $\Delta^{14}C$. The small difference between FC and RC inventories is to be attributed to the respective treatment of air-sea CO_2 fluxes. Indeed, formulation RC assumes no net air-sea CO_2 exchange at the local scale while in experiments FC these fluxes are globally, not locally, balanced. The net global CO_2 flux is zero in both cases, but spatial variability in FC slightly affects the global ^{14}C invasion rates. Experiments RP take into account the atmospheric CO_2 increase but still lag behind experiments FP. As mentioned when introducing the method consisting in transporting the ^{14}R ratio, this method implicitly assumes that air-sea CO_2 equilibrium prevails. The last term on the right-hand side of (16) fully explains the differences between the RP and FP inventories. In the eventuality of constant atmospheric ^{14}R, this term would still drive a net ^{14}C uptake by the ocean under increasing atmospheric CO_2. Not only does the atmospheric CO_2 perturbation increases exponentially after 1940, but due to the limited rate of the CO_2 exchange between the atmosphere and the ocean (typically 1 yr for a 50-m-deep column; Broecker and Peng 1974; Lynch-Stieglitz et al. 1995) and to the permanent renewal of surface waters, the ocean dissolved CO_2 increases less rapidly than atmospheric CO_2. Enhanced air-to-sea CO_2 fluxes drive a larger ^{14}C flux to the ocean in experiments FP than in experiments RP.

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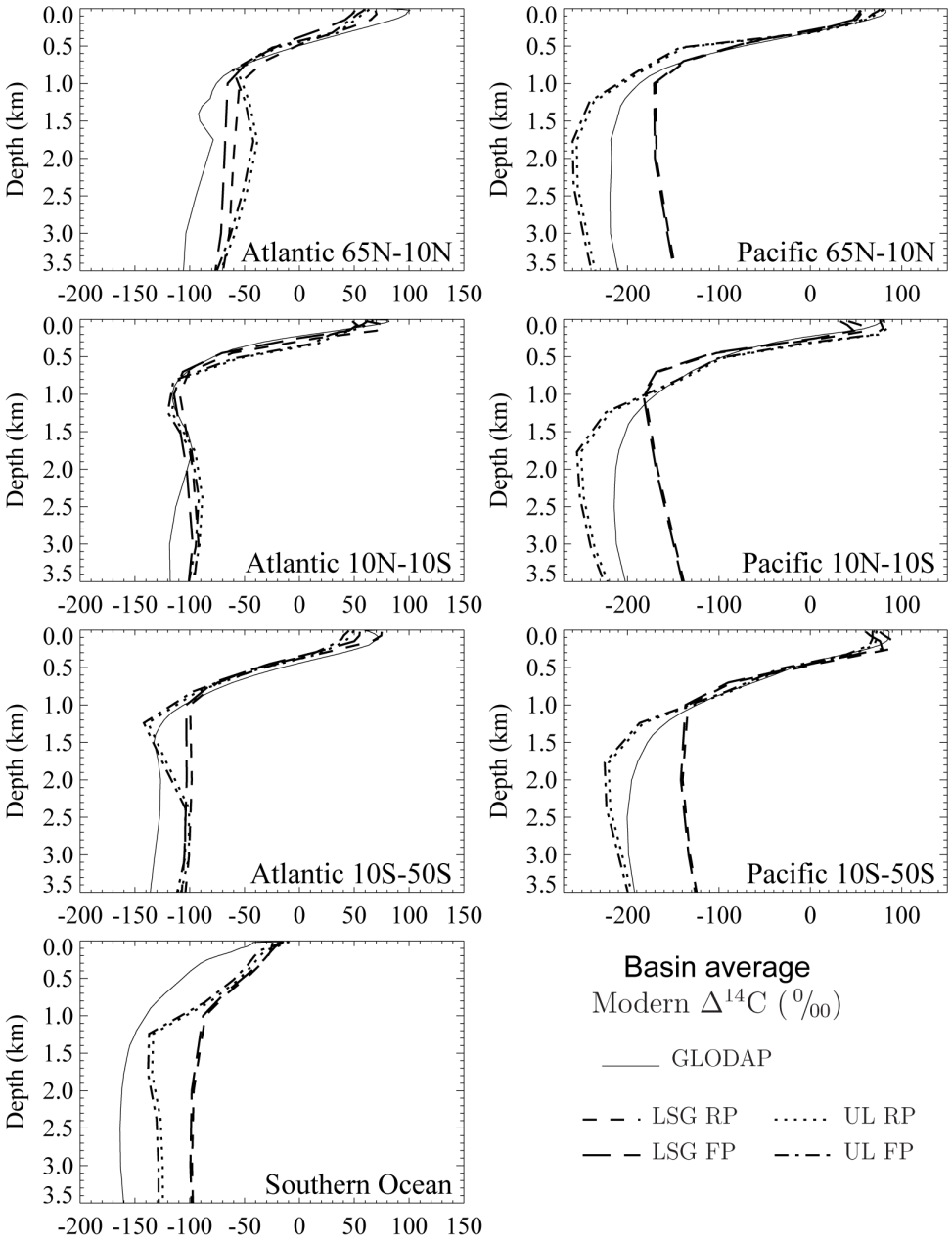


Figure 2 Mean $\Delta^{14}C$ vertical profiles in several ocean basins from GLODAP (solid lines; Key et al. 2004) and from experiments FP and RP with circulation fields from LSG-OM (LSG) and UL-OM (UL) OGCMs. The values in the north, equatorial, and southern Atlantic and Pacific oceans are reproduced, respectively, in the upper left and right 3 panels, with the bottom panel representing the Southern Ocean (south of 50°S). Model values correspond to the year 2000.

The numerical experiments allow to quantify the relative role of each process in the build up of the simulated ocean bomb ^{14}C inventories relative to 1940. $^{14}I_{1940}^{1990}$ results from 3 contributions: the atmospheric $\Delta^{14}C$ increase consecutive to the nuclear tests (exp. FC or RC) explains 80% of the FP

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inventory, the anthropogenic CO_2 increase (RP minus RC) contributes by 11% and the remaining 9% results from the enhanced air-sea CO_2 disequilibrium (FP minus RP). The contribution of the last process to the final inventory is significant; indeed, it is of the same order of magnitude as the second one. Inventory anomalies in 1990 with respect to 1975 follow the same pattern as the gross inventories ($^{14}I_{1975}^{1990}$ in Table 2). However, they exhibit larger relative differences yet. For FC and RC, $^{14}I_{1975}^{1990}$ amounts only to 73–74% of the value obtained with FP, while $^{14}I_{1975}^{1990}$ for RP represents 83–84% that of FP. It implies that over the timeframe of relevant observations (i.e. from GEOSECS to WOCE), the bomb ^{14}C inventory is extremely sensitive to the method used for computing it. Table 2 also contains the inventories obtained with a 15% lower (FPL) and a 10% larger (FPH) gas exchange velocity. The bomb ^{14}C inventory in 1990 changes by ~15% and the net CO_2 flux by ~4% over the velocity range considered here. Indeed, air-sea exchange rate exerts a dominant control on ocean ^{14}C uptake in contrast to the CO_2 uptake whose rate limiting process is the renewal time of ocean surface water with lesser control by the piston velocity (Broecker and Peng 1974; Lynch-Stieglitz et al. 1995; Ito et al. 2004).

DISCUSSION AND CONCLUSIONS

Previous works addressing the ocean bomb ^{14}C inventory with numerical models made use of different formulations of the ^{14}C isotopic ratio. Three representations are identified that yield distinct air-sea ^{14}C flux rates. In the first (method RC), the air-to-sea ^{14}C flux is simply proportional to the atmospheric $\Delta^{14}C$ signal. With the second (method RP), the actual atmospheric CO_2 concentration weights the ^{14}C flux. Eventually, with the third method (FP), which necessitates a comprehensive carbon cycle model, the CO_2 uptake by the ocean is also considered. With respect to inventories, the 3 methods lead to a larger range of results than those obtained with the circulation of 2 different OGCMs.

Inventories based on observations, on the other hand, rely on similar assumptions as method R, i.e. that DIC concentrations remain constant with time (Broecker et al. 1995; Key et al. 2004; Naegler 2009). Since reconstructed natural $\Delta^{14}C$ is constrained to match the pre-bomb observations (Rubin and Key 2002; Sweeney et al. 2007), these inventories nevertheless provide an estimate of the actual ocean $\Delta^{14}C$ increase. Therefore, data-based inventories are expected to yield similar results as method RP.

Figure 3 displays our results together with estimates from the literature. Method RC, by underestimating the ocean $\Delta^{14}C$ increase over the post-bomb times, predicts the lowest inventories. This is probably the explanation for the low estimates obtained with some previous model studies (e.g. Duffy and Caldeira 1995). The spread among model results arises from different strategies for calibrating mixing parameters and air-sea gas exchange coefficient in ocean models. The model of Joos and Bruno (1998) was tuned against the GEOSECS data, while Müller et al. (2008) optimized their model through a multi-tracer approach, in which the RP method is of use for constraining bomb ^{14}C toward the Key et al. (2004) inventory (Müller et al. 2006). Both studies afterwards applied the FP formulation in their assessment of the bomb ^{14}C inventory. In contrast to these works, I did not calibrate any of the model parameters with the help of ^{14}C data.

In order to be compared to OGCM results, inventories based on observations have to be corrected for insufficient spatial coverage as well as for lack of measurement synchronization (Naegler et al. 2006; Sweeney et al. 2007; Naegler 2009). The N06 value in Figure 3 represents the Key et al. (2004) inventory corrected for such biases (Naegler et al. 2006). The Peacock (2004) inventory also needs to be revised upward for missing ocean areas (Naegler et al. 2006). Eventually, according to Sweeney et al. (2007) their 1975 inventory may be biased low by as much as 10%. Considering

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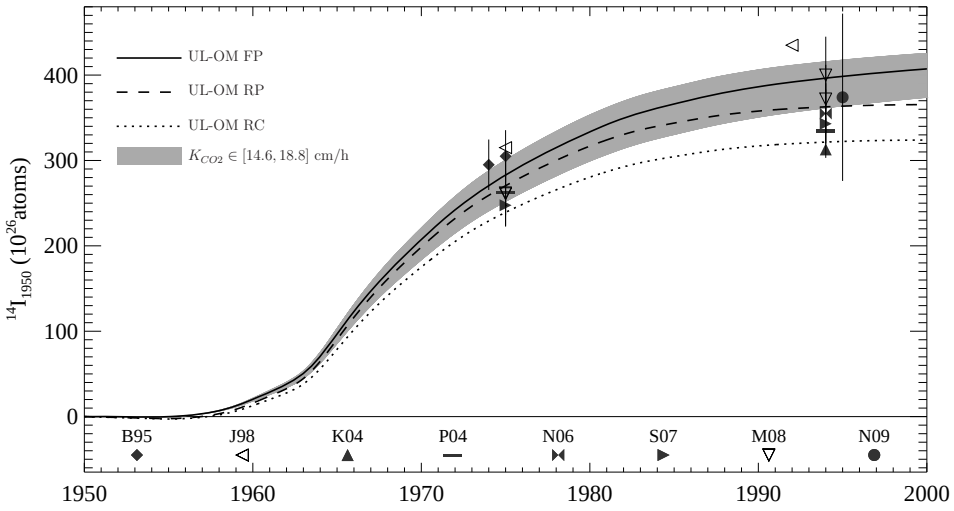


Figure 3 Ocean bomb ^{14}C inventory relative to 1950 for experiments FP, RP, and RC with UL-OM (identical curves are obtained with LSG-OM). Upper and lower bounds of the gray area correspond to experiments FPH and FPL. Estimates from the literature are also illustrated. Data (filled symbol) and model (open symbol) values as originally published by Broecker et al. (1995) B95, Joos and Bruno (1998) J98, Key et al. (2004) K04, Peacock (2004) P04, Naegler et al. (2006) N06, Sweeney et al. (2007) S07, Müller et al. (2008) M08, and Naegler (2009) N09. Where available, error bars (B95, P04, and N09) or model result range (M08) are also shown.

these facts, there is a fair agreement between the RP-modeled and data-based inventories, with the exception of the values from Broecker et al. (1995), which are probably biased high (Peacock 2004; Naegler et al. 2006).

Two important facts emerge from Figure 3. First, inventories reconstructed from field observations for which the RP predicted values are the model equivalent, underestimate the actual bomb ^{14}C inventories that are given by the FP curve. This conclusion is in agreement with the results of Naegler (2009). Second, the closeness of the results with FPL and RP in the 1990s illustrates how trying to reconcile models (FP) and data (RP) would call for a lower gas exchange rate. The results of these sensitivity studies imply that the formulation of the ^{14}C tracer in numerical models has to be adapted to the aim pursued. The global radiocarbon cycle needs to be addressed with the F (FP) method, which, by guaranteeing a full coherency between the ^{14}C and carbon air-sea exchanges, provides the actual ocean bomb ^{14}C inventory. This method is also fully relevant when investigating past or present $\Delta^{14}C$ distributions with the help of OGCMs. On the other hand, method R (RP) appears more appropriate when the purpose is to assess or calibrate upper ocean dynamics in models with the help of bomb ^{14}C distribution based on observations. This method is also relevant for the estimate of CO_2 piston velocity. For this latter purpose, the model should be assessed or calibrated without the help of bomb ^{14}C. Then one could take advantage of field estimates of the ocean bomb ^{14}C inventory to optimize the CO_2 piston velocity with method RP.

This work helps to resolve the large discrepancies between data-based estimates and model reconstructions of the bomb ^{14}C inventory in the ocean. The specification of similar time histories for the atmospheric and oceanic ^{14}C has limitations. Bomb ^{14}C inventories differ significantly, both in the 1970s and in the 1990s, between experiments where we adopted that assumption and those where we did not. Relaxing this constraint leads to bomb ^{14}C inventories that are usually larger than field estimates. Modeling studies could allow to overcome some of the difficulties inherent to field reconstruction (e.g. spatial coverage, pre-anthropogenic state). However, model results do not depend

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only on the gas exchange rate but also on the level of mixing (explicit and implicit through numerical scheme) and on the ventilation rate inherent to the model. OGCMs have deficiencies in reproducing the actual ocean ventilation rate on decadal to centennial timescales (e.g. Doney et al. 2004; Matsumoto et al. 2004), the level of vertical diffusion active in 3D OGCMs being probably too high (Duffy et al. 1997; England and Rahmstorf 1999; Oschlies 2000; Schmittner et al. 2009). Yet, this work suggests strategies allowing to gain advantage of the available data for either assessing the global carbon budget or estimating the air-sea exchange rate. Nevertheless, this problem might prove difficult to solve. Indeed, the processes involved in the ocean uptake of bomb ^{14}C (atmospheric CO_2, wind exchange, mixing, and global overturning) are characterized by widely different timescales. Available field estimates provide information on the early stages of the uptake only (~30 yr). This is possibly too short a duration to allow for a clear discrimination among the processes.

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